

Random Matrices (continued)

We now investigate the consequences of our findings on the probability distribution of eigenvalues of an arbitrary 2×2 hermitian matrix. Here again, the results would generalize to the $n \times n$ case, but we will get a sense of the issues with the 2×2 case. The basis for our approach is the fact that the trace of any matrix is simply the sum of its eigenvalues (counting multiplicity if they are repeated). So, if λ and μ are the eigenvalues of H , the eigenvalues of H^2 are clearly λ^2 and μ^2 and $\text{trace}(H^2) = \lambda^2 + \mu^2$. So, we can consider a change of variables from x_{11} , x_{12} , y_{12} , and x_{22} to λ , μ and other variables (α and β below) involved in the diagonalization of H .

From Homework Problem #2 in Class #7, we know that the most general 2×2 unitary matrix U can be written (up to a factor $e^{i\theta}$):

$$U = \begin{pmatrix} \cos\alpha & -(\sin\alpha)e^{i\beta} \\ (\sin\alpha)e^{-i\beta} & \cos\alpha \end{pmatrix} \tag{1}$$

with $\alpha \in [0, \frac{\pi}{2})$ and $\beta \in [0, 2\pi)$. However, these will depend on the order in which λ and μ are placed in the diagonal of Λ . So, we may make the convention that $\Lambda = \begin{pmatrix} \lambda & 0 \\ 0 & \mu \end{pmatrix}$ with $\lambda \geq \mu$. In that case, α and β in (1) are uniquely determined and we may proceed confidently with a change of variable. We may write (with $dH = dx_{11} dx_{22} dx_{12} dy_{12}$):

$$\int \int \int \int_{\mathcal{A}} e^{-\text{trace}(H^2)} dH = \int \int \int \int_{\mathcal{B}} e^{-(\lambda^2 + \mu^2)} |J(\lambda, \mu, \alpha, \beta)| d\lambda d\mu d\alpha d\beta \tag{2}$$

where $J(\lambda, \mu, \alpha, \beta) = \frac{\partial(x_{11}, x_{22}, x_{12}, y_{12})}{\partial(\lambda, \mu, \alpha, \beta)}$ is the usual Jacobian determinant. Calculating this determinant is what will occupy us now. Evidently, the entries of J can be found in those of $\frac{\partial H}{\partial \lambda}, \dots, \frac{\partial H}{\partial \beta}$. The only question is how to extract them in usable form. To do so, let us investigate some consequences of

$$U^* H U = \Lambda \tag{3a}$$

with

$$U^* U = I \tag{3b}$$

Differentiating (3a) with respect to λ and μ yields $U^* \frac{\partial H}{\partial \lambda} U$ and $U^* \frac{\partial H}{\partial \mu} U$ on the left and $\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$ and $\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$ on the right of (3a) respectively (since U is independent of λ and μ). Differentiating with respect to α and β requires a little more caution. Let us write (3a) as $H = U \Lambda U^*$ and let us differentiate (say with respect to α) to reach

$$\frac{\partial \mathbf{H}}{\partial \alpha} = \frac{\partial \mathbf{U}}{\partial \alpha} \Lambda \mathbf{U}^* + \mathbf{U} \Lambda \frac{\partial \mathbf{U}^*}{\partial \alpha} \quad (4a)$$

or

$$\mathbf{U}^* \frac{\partial \mathbf{H}}{\partial \alpha} \mathbf{U} = \left(\mathbf{U}^* \frac{\partial \mathbf{U}}{\partial \alpha} \right) \Lambda + \Lambda \left(\frac{\partial \mathbf{U}^*}{\partial \alpha} \mathbf{U} \right) = \mathbf{A} \Lambda + \Lambda \mathbf{A}^* \quad (4b)$$

Now, because of (3b) we have $\mathbf{U}^* \frac{\partial \mathbf{U}}{\partial \alpha} + \frac{\partial \mathbf{U}^*}{\partial \alpha} \mathbf{U} = 0 = \mathbf{A} + \mathbf{A}^*$. So:

$$\mathbf{U}^* \frac{\partial \mathbf{H}}{\partial \alpha} \mathbf{U} = \mathbf{A} \Lambda - \Lambda \mathbf{A} = (\lambda - \mu) \begin{pmatrix} 0 & s_{12}^\alpha \\ s_{21}^\alpha & 0 \end{pmatrix} \quad (4c)$$

Differentiating with β is similar to α . So we have all four matrices for $\mathbf{U}^* \frac{\partial \mathbf{H}}{\partial \xi} \mathbf{U}$ with $\xi = \lambda, \mu, \alpha,$ and β . We will now extract from those matrices the jacobian we want. We first remark that for any of our four cases of variable ξ , the entry (i, j) of $\mathbf{U}^* \frac{\partial \mathbf{H}}{\partial \xi} \mathbf{U}$ (let us denote it $(\mathbf{U}^* \frac{\partial \mathbf{H}}{\partial \xi} \mathbf{U})_{ij}$) is obtained from the same linear combination of the entries of $\frac{\partial \mathbf{H}}{\partial \xi}$. For instance:

$$\begin{aligned} (\mathbf{U}^* \frac{\partial \mathbf{H}}{\partial \xi} \mathbf{U})_{ij} &= l_{11}^{ij} \frac{\partial x_{11}}{\partial \xi} + l_{22}^{ij} \frac{\partial x_{22}}{\partial \xi} + l_{12}^{ij} \frac{\partial x_{12}}{\partial \xi} + l_{21}^{ij} \frac{\partial y_{12}}{\partial \xi} \\ &= \left(\frac{\partial x_{11}}{\partial \xi}, \frac{\partial x_{22}}{\partial \xi}, \frac{\partial x_{12}}{\partial \xi}, \frac{\partial y_{12}}{\partial \xi} \right) (l_{11}^{ij}, l_{22}^{ij}, l_{12}^{ij}, l_{21}^{ij})^T \end{aligned} \quad (5)$$

So, if we let $\mathbf{L}_{ij} = (l_{11}^{ij}, l_{22}^{ij}, l_{12}^{ij}, l_{21}^{ij})^T$ and let \mathbf{L} be the determinant with columns \mathbf{L}_{ij} as follows: $\mathbf{L} = |\mathbf{L}_{11}, \mathbf{L}_{22}, \mathbf{L}_{12}, \mathbf{L}_{21}|$, we can easily write the *product* determinant

$$\mathbf{J}\mathbf{L} = \begin{vmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & s_{12}^\alpha (\lambda - \mu) & s_{21}^\alpha (\lambda - \mu) \\ 0 & 0 & s_{12}^\beta (\lambda - \mu) & s_{21}^\beta (\lambda - \mu) \end{vmatrix} = (\lambda - \mu)^2 \begin{vmatrix} s_{12}^\alpha & s_{21}^\alpha \\ s_{12}^\beta & s_{21}^\beta \end{vmatrix} = \mathbf{S}(\lambda - \mu)^2$$

Now, clearly, all the elements of \mathbf{L} and \mathbf{S} depend solely on α and β (not on λ and μ). So, we can write $|\mathbf{J}| = \frac{|\mathbf{S}|}{|\mathbf{L}|} (\lambda - \mu)^2 = g(\alpha, \beta) (\lambda - \mu)^2$. Replacing in (2) yields

$$\begin{aligned} &\int \int \int \int_{\mathcal{B}} e^{-(\lambda^2 + \mu^2)} |\mathbf{J}(\lambda, \mu, \alpha, \beta)| d\lambda d\mu d\alpha d\beta \\ &= c \int \int_{\mathcal{B}} e^{-(\lambda^2 + \mu^2)} (\lambda - \mu)^2 d\lambda d\mu \end{aligned} \quad (6)$$

where c is a normalizing constant that includes the integration of $g(\alpha, \beta)$.

Formula (6) shows the form of the probability density function of the eigenvalues λ and μ for a 2×2 hermitian matrix. The entire above construction can be generalized to the $n \times n$ case and reads:

$$c_n \int \int_{\mathcal{B}} e^{-\sum_{i=1}^n \lambda_i^2} \prod_{j < i} (\lambda_i - \lambda_j)^2 d\lambda_1 \dots d\lambda_n \quad (7)$$

The square factor characterizes the hermitian case. Other cases involve other powers of the difference of eigenvalues.

Homework

1. In (6), given λ , find μ where the probability density is maximum. Write an expression for the mean value of μ (do not estimate it).